

# Neutron Interactions and Dosimetry

## Chapter 16

F.A. Attix, Introduction to Radiological Physics and Radiation Dosimetry

## Outline

- Neutron dosimetry
  - Thermal neutrons
  - Intermediate-energy neutrons
  - Fast neutrons
- Sources of neutrons
- Mixed field dosimetry, paired dosimeters
- Rem meters

## Introduction

- Consider neutron interactions with the majority tissue elements H, O, C, and N, and the resulting absorbed dose
- Because of the short ranges of the secondary charged particles that are produced in such interactions, CPE is usually well approximated
- Since no bremsstrahlung x-rays are generated, the absorbed dose can be assumed to be equal to the kerma at any point in neutron fields at least up to an energy  $E \sim 20$  MeV

## Tissue composition

Element	ICRU (1977)	ICRP Skeletal	Whole Body
	Muscle	Muscle	
H	10.2	10.06	10.5
O	72.9	75.48	67.7
C	12.3	10.78	18.7
N	3.5	2.77	3.1
Other	1.1	0.91	—
Total	100.0	100.00	100.0

\*Data are given in percent by weight.

- The ICRU composition for muscle has been assumed in most cases for neutron-dose calculations, lumping the 1.1% of “other” minor elements together with oxygen to make a simple four-element (H, O, C, N) composition

## Neutron kinetic energy

- Neutron fields are divided into three categories based on their kinetic energy:
  - Thermal ( $E < 0.5$  eV)
  - Intermediate-energy ( $0.5 \text{ eV} < E < 10$  keV)
  - Fast ( $E > 10$  keV)
- Differ by their primary interactions in tissue and resulting biological effects

## Neutron kinetic energy

- Thermal neutrons, by definition, have the most probable kinetic energy  $E = kT = 0.025 \text{ eV}$  at  $T = 20^\circ\text{C}$
- Neutrons up to 0.5eV are considered “thermal” due to simplicity of experimental test after they emerge from moderator material
- Cadmium ratio test:
  - Gold foil can be activated through  $^{197}\text{Au}(n,\gamma)^{198}\text{Au}$  interaction
  - Addition of 1mm thick Cd filter, which absorbs all neutrons below 0.5eV, tests for presence of those neutrons

## Neutron kinetic energy

- For  $E < 10$  keV the dose is mainly due to  $\gamma$ -rays resulting from thermal-neutron capture in hydrogen,  ${}^1\text{H}(n,\gamma){}^2\text{H}$
- For  $E > 10$  keV the dose is dominated by the contribution of recoil protons resulting from elastic scattering of hydrogen nuclei
  - Resulting biological effect (and the corresponding neutron quality factor) rises steeply

## Kerma calculations

- For a mono-energetic neutron beam having fluence  $\Phi$  ( $\text{cm}^{-2}$ ), the kerma that results from a neutron at a point in a medium is

$$K = 1.602 \times 10^{-8} \Phi \sigma N_t m^{-1} E_{tr}$$

where  $\sigma$  is the interaction cross section in  $\text{cm}^2/(\text{target atom})$ ,  $N_t$  is the number of target atoms in the irradiated sample,  $m$  is the sample mass in grams, and  $E_{tr}$  is the total kinetic energy (MeV) given to charged particles per interaction

## Kerma calculations

- The product of  $(1.602 \times 10^{-8} \sigma N_t m^{-1} E_{tr})$  is equal to the kerma factor  $F_n$  in  $\text{rad cm}^2/n$
- Thus the equation reduces to

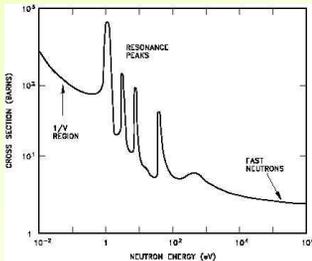
$$D = K = \Phi F_n$$

- CPE condition is usually well-approximated for neutron interactions in tissue

## Kerma calculations

- $F_n$  values are tabulated in Appendix F
- $F_n$  is not generally a smooth function of  $Z$  and  $E$ , unlike the case of photon interaction coefficients
- Interpolation vs.  $Z$  cannot be employed to obtain values of  $F_n$  for elements for which data are not listed, and interpolation vs.  $E$  is feasible only within energy regions where resonance peaks are absent

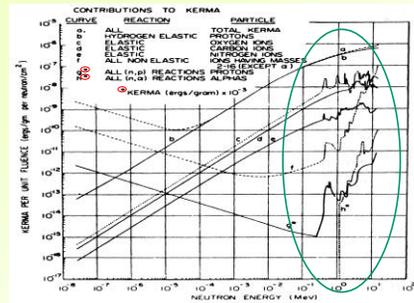
## Typical neutron absorption cross section



Resonance peaks occur when the binding energy of a neutron plus the KE of the neutron are exactly equal to the amount required to raise a compound nucleus from its ground state to one of quantum levels

Image from [http://nuclearpowertraining.epub.com/h1019v1/css/h1019v1\\_113.htm](http://nuclearpowertraining.epub.com/h1019v1/css/h1019v1_113.htm)

## Kerma calculations



Kerma per unit fluence due to various interaction in a small mass of tissue

## Thermal-neutron interactions in tissue

- Two important interactions of thermal neutrons with tissue:
  - neutron capture by nitrogen,  $^{14}\text{N}(n,p)^{14}\text{C}$
  - neutron capture by hydrogen,  $^1\text{H}(n,\gamma)^2\text{H}$
- Thermal neutrons have a larger probability of capture by hydrogen atoms in muscle, because there are 41 times more H atoms than N atoms in tissue

## Thermal-neutron interactions in tissue

- The nitrogen interaction releases a kinetic energy of  $E_{tr} = 0.62$  MeV that is shared by the proton (0.58 MeV) and the recoiling nucleus (0.04 MeV)
- CPE exists since the range of protons in tissue  $\sim 10\mu\text{m}$
- Based on known values for  $\sigma$  and  $N/m$ , kerma factor
 
$$F_n = 2.74 \times 10^{-11} [\text{rad cm}^2 / n]$$
- Dose deposited  $^{CPE}$ 

$$D = 2.74 \times 10^{-11} \Phi [\text{rad}]$$

## Thermal-neutron interactions in tissue

- In the hydrogen interaction the energy given to  $\gamma$ -rays per unit mass and per unit fluence of thermal neutrons can be obtained similarly, but replacing  $E_{tr}$  by  $E_\gamma = 2.2$  MeV (the  $\gamma$ -ray energy released in each neutron capture)
- The energy available:  $R_\gamma / \Phi m = 7.13 \times 10^{-10} [\text{rad cm}^2 / n]$
- These  $\gamma$ -rays must interact and transfer their energy to charged particles to produce kerma
- Contribution to dose depends on proximity to RE

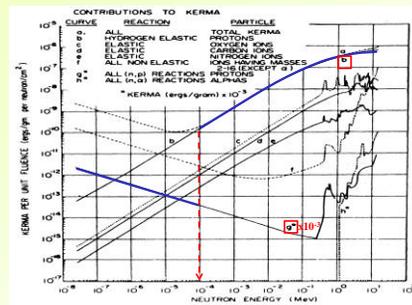
## Thermal-neutron interactions in tissue

- In the center of a 1-cm diameter sphere of tissue the kerma contributions from the  $(n, p)$  and  $(n, \gamma)$  processes are comparable in size
- In a large tissue mass (radius  $> 5$  times the  $\gamma$ -ray MFP) where radiation equilibrium is approximated, the kerma due indirectly to the  $(n, \gamma)$  process is 26 times that of the nitrogen  $(n, p)$  interaction
- The human body is intermediate in size, but large enough so the  $^1\text{H}(n, \gamma)^2\text{H}$  process dominates in kerma (and dose) production

## Interactions by intermediate and fast neutrons in tissue

- It is dominated below  $10^{-4}$  MeV (100eV) by  $(n,p)$  reactions, mostly in nitrogen, represented by curve *g*
- Above  $10^{-4}$  MeV elastic scattering of hydrogen nuclei (curve *b*) contributes nearly all of the kerma

## Interactions by neutrons in tissue



## Interaction by intermediate and fast neutrons in tissue

- The average energy transferred by elastic scattering to a nucleus is closely approximated (i.e., assuming isotropic scattering in the center-of-mass system) by

$$\bar{E}_{tr} = E \frac{2M_a M_n}{(M_a + M_n)^2}$$

where  $E$  = neutron energy,

$M_a$  = mass of target nucleus,

$M_n$  = neutron mass

## Interaction by intermediate and fast neutrons in tissue

- Values for  $\bar{E}_{tr}$  for different tissue atoms:
  - $E/2$  for hydrogen recoils (changes from 0 when proton is recoiling at  $90^\circ$  to  $E_{tr}$  for protons recoiling straight ahead)
  - $0.142E$  for C atoms
  - $0.124E$  for N
  - $0.083E$  for O

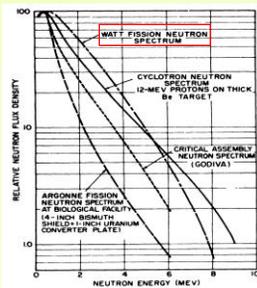
## Neutron sources

- Neutrons can be released through processes such as fission, spallation, or neutron stripping
- Fission: splitting atoms in a nuclear reactor
  - $n + {}^{235}\text{U} = n + n + \text{fragments}$  – one n may go back into chain reaction, the other is available
- Spallation: bombarding heavy metal atoms with energetic protons
  - $p + \text{heavy nucleus} = X \cdot n + \text{fragments}$
- Stripping: bombarding light metal atoms with protons,  $\alpha$ -particles, etc.

## Neutron sources

- Most widely available neutron sources:
  - Nuclear fission reactors
  - Accelerators
  - Radioactive sources

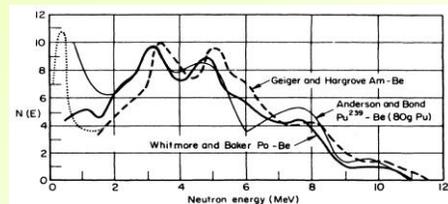
## Neutron sources



Fission neutron spectra from reactors

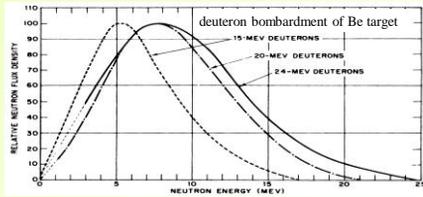
- Fission-like spectra have average neutron energies  $\sim 2\text{MeV}$
- From beam ports of nuclear reactors
- Watt spectrum is an idealized shape for unmoderated reactor neutrons

## Neutron sources



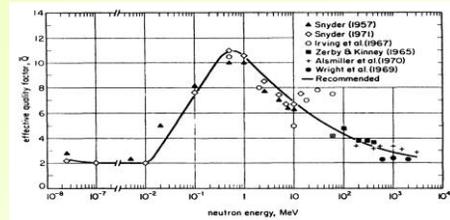
- Several types of  $\text{Be}(\alpha, n)$  radioactive sources are in common use, employing  ${}^{210}\text{Po}$ ,  ${}^{239}\text{Pu}$ ,  ${}^{241}\text{Am}$ , or  ${}^{226}\text{Ra}$  as the emitter
- Neutron yields are on the order of 1 neutron per  $10^4$   $\alpha$ -particles; average energies  $\cong 4$  MeV

## Neutron sources



- Cyclotrons are used to produce neutron beams by accelerating protons or deuterons into various targets, most commonly Be
- The neutron spectrum extends from 0 to energies  $E_{\max}$  somewhat above the deuteron energy, and have  $E_{\text{avg}} \sim 0.4E_{\max}$
- Tissue dose rates  $\sim 10 \text{ rad/min}$  at 1m in tissue,  $\gamma$ -ray background  $\sim$  few %

## Neutron Quality Factor



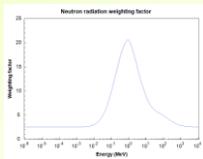
- For purposes of neutron radiation protection the dose equivalent  $H = DQ$ , where  $Q$  is the quality factor, which depends on neutron energy
- The quality factor for all  $\gamma$ -rays is taken to be unity for purposes of combining neutron and  $\gamma$ -ray dose equivalents

## Radiation Weighting Factor

- New approach uses *radiation and tissue weighting factors* (updated in ICRP 103 report, 2007); effective dose

$$E = w_T H_T = w_T w_R^* D_{T,R}$$

- For neutrons weighting factor is a continuous function of energy; represented as a graph or in functional form



$$w_R = \begin{cases} 2.5 + 18.2e^{-\ln(E_n/0.046)}^2 & E_n < 1 \text{ MeV} \\ 5.0 + 17.0e^{-\ln(E_n/0.046)} & 1 \text{ MeV} \leq E_n \leq 50 \text{ MeV} \\ 2.5 + 3.25e^{-\ln(E_n/0.046)}^2 & E_n > 50 \text{ MeV} \end{cases}$$

## Mixed-Field Dosimetry $n + \gamma$

- Neutrons and  $\gamma$ -rays are both indirectly ionizing radiations that are attenuated exponentially in passing through matter
- Each is capable of generating secondary fields of the other radiation:
  - $(\gamma, n)$  reactions are only significant for high-energy  $\gamma$ -rays ( $\geq 10 \text{ MeV}$ , there is a threshold defined by reaction)
  - $(n, \gamma)$  reactions can proceed at all neutron energies and are especially important in thermal-neutron capture, as discussed for  ${}^1\text{H}(n, \gamma){}^2\text{H}$

## Mixed-Field Dosimetry $n + \gamma$

- Neutron fields are normally “contaminated” by secondary  $\gamma$ -rays
- Since neutrons generally have more biological effectiveness per unit of absorbed dose than  $\gamma$ -rays, it is desirable to account for  $\gamma$  and  $n$  components separately
- It is especially important in the case of neutron dosimeters to specify the reference material to which the dose reading is supposed to refer

## Mixed-Field Dosimetry $n + \gamma$

- Three general categories of dosimeters for  $n + \gamma$  applications:
  1. Neutron dosimeters that are relatively insensitive to  $\gamma$  rays
  2.  $\gamma$ -ray dosimeters that are relatively insensitive to neutrons
  3. Dosimeters that are comparably sensitive to both radiations

## Mixed-Field Dosimetry $n + \gamma$

- While water is a very close substitute of muscle tissue in photon dosimetry, it is not as close a substitute for neutrons
- Water is 1/9 hydrogen by weight; muscle is 1/10 hydrogen
- Water contains no nitrogen, and hence can have no  $^{14}\text{N}(n,p)^{14}\text{C}$  reactions by thermal neutrons
- For 1MeV photon  $(\mu_{\text{en}}/\rho)_{\text{muscle}}=0.99 (\mu_{\text{en}}/\rho)_{\text{water}}$  ,  
for 1MeV neutron  $(F_n)_{\text{muscle}}=0.91(F_n)_{\text{water}}$

## Mixed-Field Dosimetry $n + \gamma$

- The response of a dosimeter to a mixed field of neutrons and  $\gamma$ -rays

$$Q_{n,\gamma} = AD_\gamma + BD_n$$

or alternatively as

$$\frac{Q_{n,\gamma}}{A} = D_\gamma + \frac{B}{A}D_n$$

Here  $A$  = response per unit of absorbed dose in tissue for  $\gamma$ -rays,  $B$  = response per unit of absorbed dose in tissue for neutrons,  $D_\gamma$  and  $D_n$  are absorbed doses in tissue correspondingly due to  $\gamma$ -rays and neutrons

## Mixed-Field Dosimetry $n + \gamma$

- By convention the absorbed dose referred to in these terms is assumed to be that under CPE conditions in a small imaginary sphere of muscle tissue, centered at the dosimeter midpoint with the dosimeter absent
- Most commonly this tissue sphere is taken to be just large enough (0.52-g/cm<sup>2</sup> radius) to produce CPE at its center in a  $^{60}\text{Co}$  beam

## Mixed-Field Dosimetry $n + \gamma$

- An approach of employing two dosimeters (paired dosimeters) having different values of  $A/B$  can be used to simultaneously obtain  $D_\gamma$  and  $D_n$
- The best dosimeter pair is a TE-plastic ion chamber containing TE gas (for which  $B/A \cong 1$ ) to measure the total  $n + \gamma$  dose, and a nonhydrogenous dosimeter having as little neutron sensitivity as possible to measure the  $\gamma$  dose
  - Ideally this paired dosimeter should measure *only*  $\gamma$ -rays
- Closer values of  $A/B$  in a pair decrease the accuracy

## Mixed-Field Dosimetry $n + \gamma$

- Often dosimeters insensitive to  $\gamma$ -rays are used in mixed fields to evaluate neutron dose only ( $A \ll B$ )
- ★ – Activation of metal foils ( $A \cong 0$ )
  - Fission foils ( $A = 0$ )
  - Etchable plastic foils ( $A \cong 0$ )
  - Damage to silicon diodes ( $A \cong 0$ )
  - Hurst proportional counter ( $A \cong 0$ )
- ★ – Rem meters
  - Long counters
  - Bubble detectors

## Activation of metal foils

- Since most of radio-activation by photonuclear reactions can occur only above the energy range of  $\gamma$ -radiation ( $>10\text{MeV}$ ), metal foils are only activated by the neutrons in the mixed field
- The resulting activity of a foil is measured by counting  $\gamma$ -rays emitted (G-M counter will work)
- Some activated foils are  $\beta$ -emitters (e.g., in  $^{32}\text{S}(n,p)^{32}\text{P}$  reaction,  $^{32}\text{P}$  is a  $\beta$ -emitter)

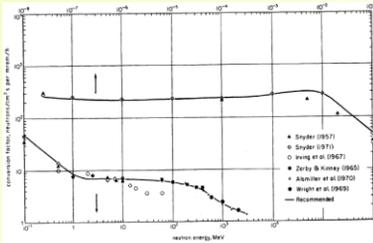
## Activation of metal foils

- Thermal neutrons have a fixed activation cross section  $\sigma(E)=\sigma$
- Fast neutrons have activation threshold, below which  $\sigma(E)=0$
- Since different materials have different threshold, a set of foils with different thresholds allows determination of the neutron spectrum

## Rem meters

- Instruments designated to measure dose equivalent  $H$ , evaluated at the depth in the body where it reaches a maximum for each incident neutron energy
- The ICRP provided the fluence-to-maximum dose-equivalent conversion factor  $d(E)$  (units of mrem/(n cm<sup>2</sup>)) to the neutron energy incident on the body
- Dose equivalent  $H = \int_{E=0}^{E_{max}} \Phi'(E)d(E)dE$ , here  $\Phi'(E)$  is the energy spectrum of incident neutrons

## Rem meters



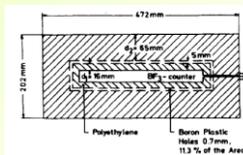
The maximum value of the dose equivalent,  $[3600d(E)]^{-1}$  in the body (tissue slab in 30-cm thickness) vs energy of neutrons

## Rem meters

- Instruments are based on materials having high capture cross section for neutrons
- Most rem meters employ boron in form of BF<sub>3</sub> gas (very toxic and corrosive); newer instruments use helium gas
- Reactions:
  - $^{10}\text{B} + n \rightarrow ^7\text{Li} + \alpha + 2790 \text{ keV}$
  - $^3\text{He} + n \rightarrow ^3\text{H} + ^1\text{H} + 765 \text{ keV}$

## Rem meters

- For both  $^{10}\text{B}$  and  $^3\text{H}$  the cross section  $\sigma \sim 1/\sqrt{E} \sim 1/v$
- Require neutron moderator, typically polyethylene
- Response is proportional to the absorbed dose and the neutron quality



Sectional view of Andersson-Braun rem counter

## Summary

- Neutron dosimetry approaches differ from those of photon dosimetry
- Water is not the best tissue-mimicking phantom anymore
- Quality factor for neutrons is energy dependent
- Dosimetry involves measurement of mixed n+ $\gamma$  fields